

# Waveguide Composed of Pinhole Array

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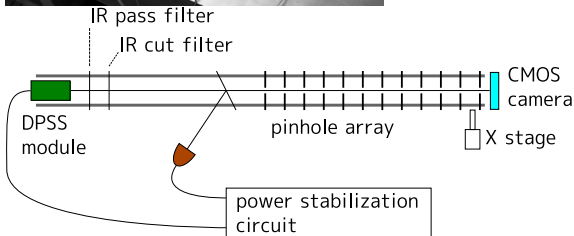
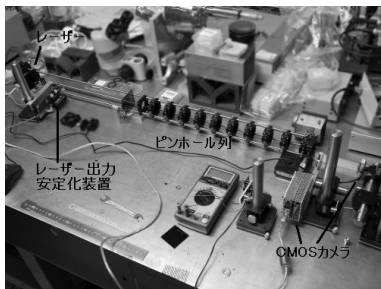
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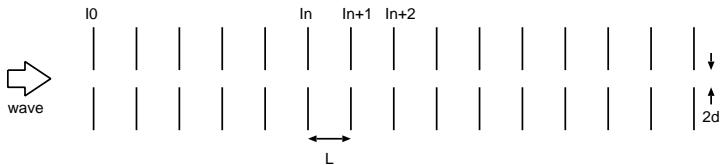
# abstract

A waveguide composed of an array of pinholes is proposed. Intensity loss at each pinhole is shown to be proportional to  $L^{1.5}$  where  $L$  is the distance between adjacent pinholes. Thus by increasing the number of pinholes 4 times, the loss for unit length becomes a half. Experimental demonstration of such waveguide is also given.

# guiding of wave (light, atoms,...) with pinholes/slits



$\lambda = 1064\text{nm}, 532\text{nm}$



$$I_{n+1} = \left\{ 1 - \beta \left( \frac{\lambda L}{d^2} \right)^{\frac{3}{2}} \right\} I_n$$

By reducing the pinhole/slit spacing  $1/4$  the loss per unit length becomes a half!

## Propagation through the pinhole array

Propagation process can be divided into two parts:

- 1 Free propagation between two neighboring slits

$$\mathbf{T}_L = \int_{-\infty}^{\infty} dk |k\rangle e^{ik_z L} \langle k| = e^{i\phi_0} \int_{-\infty}^{\infty} dk |k\rangle e^{-i\alpha k^2} \langle k|$$

$$k_0^2 = k_z^2 + k^2 \rightarrow k_z \sim k_0 - \frac{k^2}{2k_0} \rightarrow \alpha = \frac{L}{2k_0}$$

- 2 Masking of wavefunction when the wave pass through a slit

$$(\mathbf{T}_M \psi)(\mathbf{x}) = \begin{cases} \psi(\mathbf{x}) & |\mathbf{x}| < d \\ 0 & \text{otherwise} \end{cases}$$

- 3  $\mathbf{T} \equiv \mathbf{T}_M \mathbf{T}_L$

- 4 Propagation through many slits:  $\mathbf{T}^n$

- 5 Find eigenvalue and eigenstate of  $\mathbf{T}$

## Wavefunction just after a slit

- Takes nonzero value only inside the opening
- Can be expanded with sin/cos function that are 0 on the slit boundary:

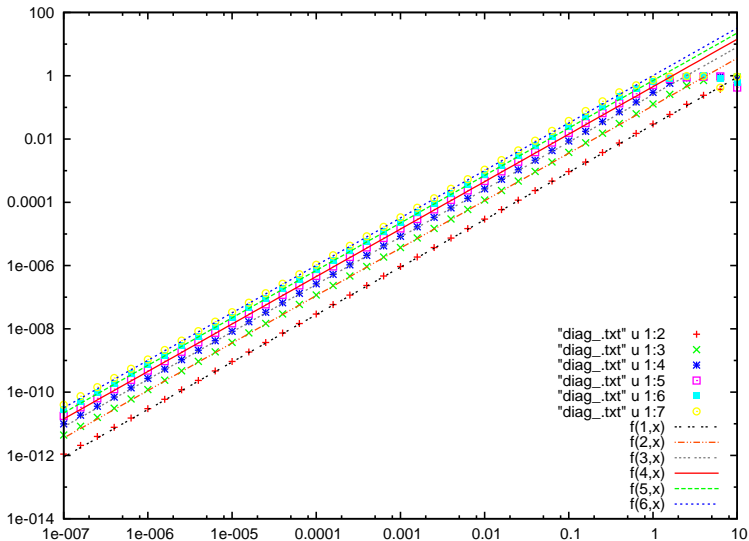
$$\psi_n(x) = \begin{cases} c_0 \cos k_n x & (n : \text{even}) \\ c_0 \sin k_n x & (n : \text{odd}) \\ 0 & (\text{otherwise}) \end{cases} \quad (-d \leq x \leq d)$$

$$k_n = \frac{n+1}{2d} \pi \quad (n = 0, 1, 2, \dots)$$

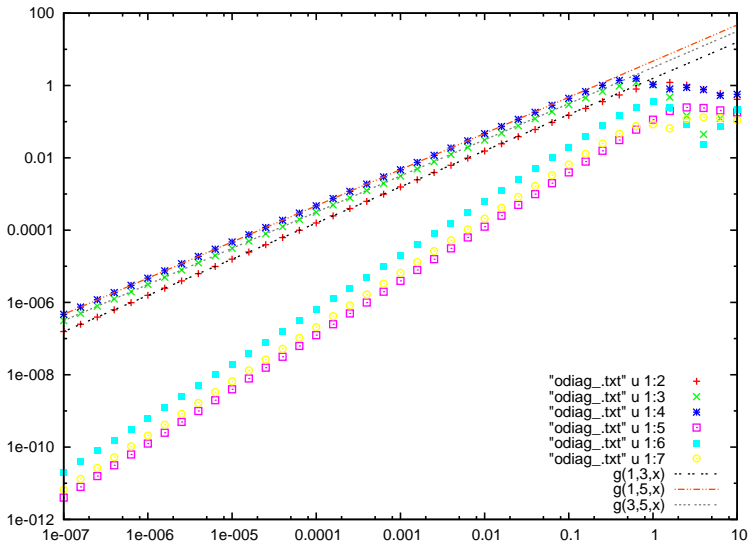
$$c_0 = \frac{1}{\sqrt{d}}: \text{normalization constant}$$

( $\psi_n$  are eigenfunctions of “Particle in a Box”)

$$1 - |T_{nn}|$$



# off diagonal elements v.s. difference of the diagonal elements of $\mathbf{T}$





asymptotic form of the operator  $\mathbf{T}$ 

$$\begin{aligned}\mathbf{T}_{mn} = & \delta_{mn} \exp\left(-i \frac{(m+1)^2 \pi^2}{8} \frac{L}{kd^2}\right) \\ & + \frac{1+(-1)^{m+n}}{2} (m+1)(n+1) \frac{\pi^{\frac{3}{2}}}{12} \left(\frac{L}{kd^2}\right)^{\frac{3}{2}} (-1-i) \\ & + \mathcal{O}\left(\left(\frac{L}{kd^2}\right)^2\right)\end{aligned}$$

amplitude loss per single unit of slit (length  $L$ ):

$$1 - |\mathbf{T}_{mm}| = (m+1)^2 \frac{\pi^{\frac{3}{2}}}{12} \left(\frac{L}{kd^2}\right)^{\frac{3}{2}} + \mathcal{O}\left(\left(\frac{L}{kd^2}\right)^2\right)$$

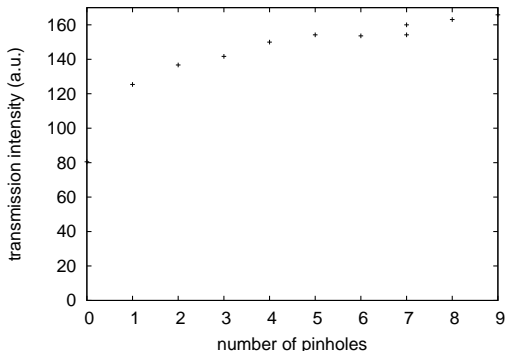
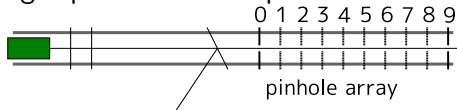
# eigenstates & eigenvalues of $\mathbf{T}$

- difference of the diagonal elements:  $\mathbf{O}\left(\frac{\lambda L}{d^2}\right)$
- off-diagonal elements:  $\mathbf{O}\left(\frac{\lambda L}{d^2}\right)^{1.5}$

When  $\frac{\lambda L}{d^2} \rightarrow 0$

- $\psi_n$  become eigenstates
- $\mathbf{T}_{nn}$  become eigenvalues

# light power after the pinhole no.9



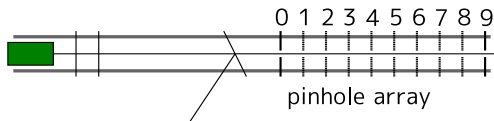
$$\lambda = 532\text{nm}$$

$$\phi = 0.5\text{mm}$$

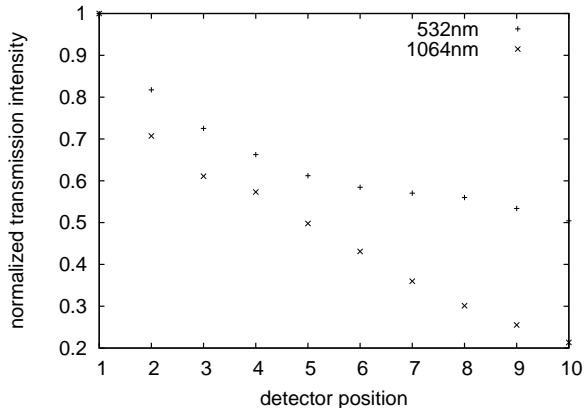
$$L = 45\text{mm}$$

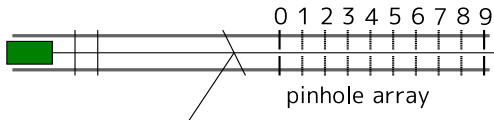
- 0: pinhole no.9 only,
- 1: pinhole no.9 + no.0,
- 2: pinhole no.9 + no.0 + no.1,
- ...

160 in the vertical axis corresponds to about 50% of light power transmitted through pinhole no.0.

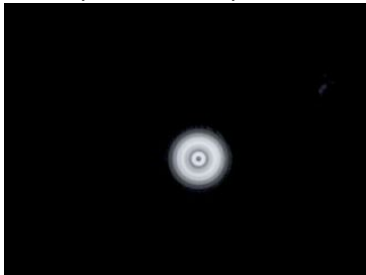


light power after each pinhole



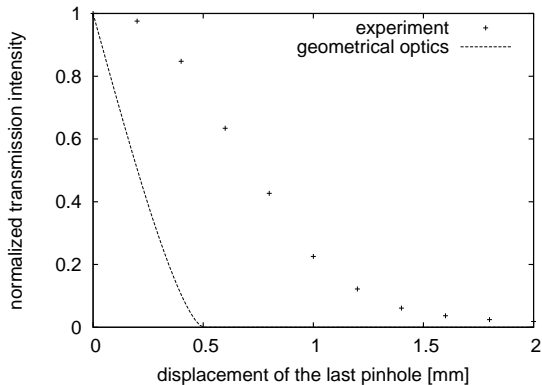
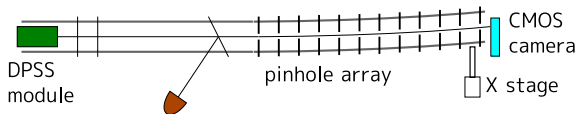


beam profile after pinhole no.0 and no.3



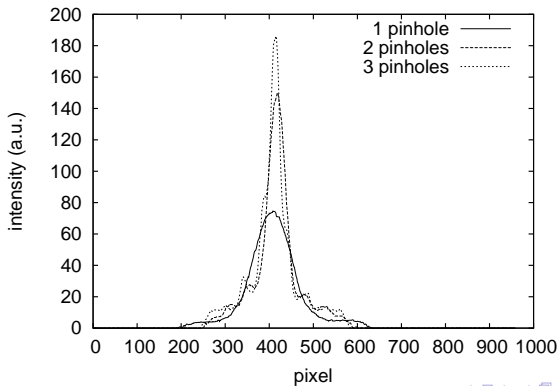
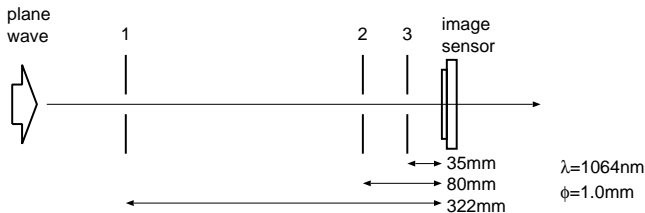
$$\lambda = 532\text{nm}, \phi = 0.5\text{mm}, L = 45\text{mm}$$

## beam deflection



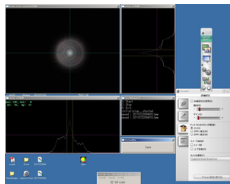
$$\lambda = 532\text{nm}, \phi = 0.5\text{mm}, L = 45\text{mm}$$

# beam focusing



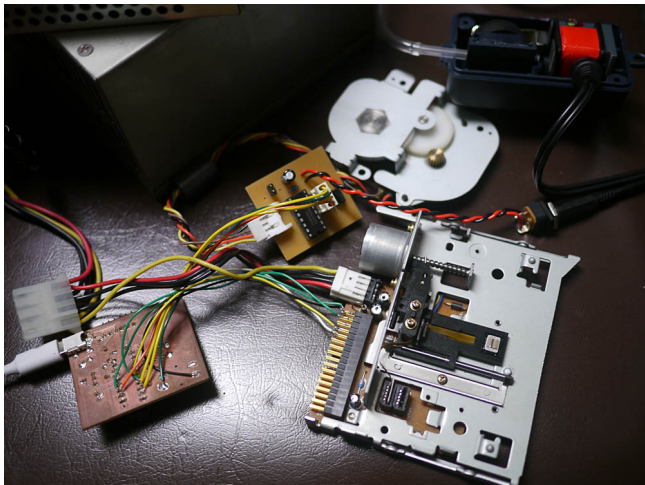
## features of this methods application

- Any kind of wave can be handled
  - electro-magnetic wave of special wavelength, matter wave, etc.
- possible to bend
- wave travels through a free space
  - guide atoms by the light guided by the pinhole array
- manipulation of the transverse mode
  
- **FREE** beam profiling program for WDM cameras (USB cameras) is available at: <http://m.ils.uec.ac.jp/sbpw/>





# pinhole aligning ROBOT (under construction)



number of pinholes:  $\sim 100$

pinhole spacing:  $\gtrsim 2\text{mm}$

naive estimate of the loss  
coefficient

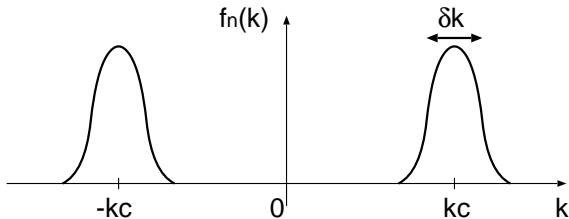
- projection of  $\mathbf{T}\psi$  onto  $\psi$

$$\mathbf{S} = \langle \psi | \mathbf{T} | \psi \rangle = \langle \psi | \mathbf{T}_L | \psi \rangle = \int d\mathbf{k} |\langle \psi | \mathbf{k} \rangle|^2 e^{-i\alpha \mathbf{k}^2}$$

- $\mathbf{1} - |\mathbf{S}_n|^2 = \mathbf{1} - |\langle \psi_n | \mathbf{T} | \psi_n \rangle|^2$  gives the intensity loss coefficient for the  $n$ th mode when  $\frac{\lambda L}{d^2} \ll 1$
- property of  $\mathbf{f}(\mathbf{k}) \equiv |\langle \psi | \mathbf{k} \rangle|^2$ :

$$\int_{-\infty}^{\infty} d\mathbf{k} \mathbf{f}(\mathbf{k}) = 1, \quad \mathbf{f}(\mathbf{k}) \geq 0$$

$$f_n(\mathbf{k}) \equiv |\langle \psi_n | \mathbf{k} \rangle|^2$$



$$k_c = \frac{n+1}{2d} \pi, \quad \delta k \sim \frac{\pi}{2d}$$

$$k_c \equiv \int_0^{\infty} dk \, k \{f(k) + f(-k)\}$$

$$\delta k \equiv \left( \int_0^{\infty} dk \, (k - k_c)^2 \{f(k) + f(-k)\} \right)^{\frac{1}{2}}$$

$$\begin{aligned}
 & \mathbf{S} \\
 = & \langle \psi | \mathbf{T} | \psi \rangle \\
 = & \int_{-\infty}^{\infty} dk f(k) e^{-i\alpha k^2} \\
 = & e^{-i\alpha k_c^2} \int_0^{\infty} dk \{f(k) + f(-k)\} e^{-i\alpha(2k_c \Delta k + \Delta k^2)} \\
 \sim & e^{-i\alpha k_c^2} \int_0^{\infty} dk g(k) \{1 - i\alpha(2k_c \Delta k + \Delta k^2) - 2\alpha^2 k_c^2 \Delta k^2\} \\
 = & e^{-i\alpha k_c^2} \{1 + (-i\alpha - 2\alpha^2 k_c^2) \delta k^2\}
 \end{aligned}$$

$$|S|^2 \sim 1 - 4\alpha^2 k_c^2 \delta k^2 = 1 - \frac{k_c^2 \delta k^2}{k_0^2} L^2$$

$$|S_n|^2 \sim 1 - \frac{(n+1)^2 \pi^4}{16k_0^2 d^4} L^2 = 1 - \frac{(n+1)^2 \pi^2}{64} \left(\frac{\lambda L}{d^2}\right)^2$$

wrong result!

$$g(k) = f(k) + f(-k), \Delta k \equiv k - k_c.$$